

# On Fast Algorithms for Computing Spatial Distance Histograms (SDH)

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## Abstract

Particle simulation has become an important research tool in many scientific and engineering fields. Data generated by such simulations impose great challenges to database storage and query processing. One of the queries against particle simulation data, the spatial distance histogram (SDH) query, is the building block of many high-level analytics, and requires quadratic time to compute using a straightforward algorithm. Previous work has developed efficient algorithms that compute exact SDHs with time complexity  $O(N^{\frac{3}{2}})$  for two-dimensional data, and  $O(N^{\frac{5}{3}})$  for three-dimensional data. While beating the naive solution, such algorithms are still not practical in processing SDH queries against large-scale simulation data. In this paper, we take a different path to tackle this problem by focusing on approximate algorithms with provable error bounds. We first present a solution derived from the aforementioned exact SDH algorithm, and this solution has running time that is unrelated to the input size  $N$ . While an error bound can be easily identified, experimental results show that the accuracy of such an algorithm is significantly higher than what is given by such a (loose) bound. To study the difference between the experimental results and the theoretical bound, we develop a mathematical model to analyze the mechanism that leads to errors in the basic approximate algorithm. Our model provides insights on how the algorithm can be improved to achieve higher accuracy and efficiency. Such insights give rise to a new approximate algorithm with improved time/accuracy tradeoff. Experimental results confirm our analysis well.

## Index Terms

molecular simulation, particle simulation, spatial distance histogram, radial distribution functions, quad-tree, scientific databases

## I. INTRODUCTION

Many scientific fields have undergone a transition to data/computation intensive science, as the result of automated experimental equipments and computer simulations. In recent years, much progress has been made in building data management tools suitable for processing scientific data [1]–[5]. Scientific data imposes great challenges to the design of database management systems that are traditionally optimized toward handling business applications. First, scientific data often come in large volumes, this requires us to rethink the storage, retrieval, and replication techniques in current DBMSs. Second, user accesses to scientific databases are focused on complex high-level analytics and reasoning that go beyond simple aggregate queries. While many types of domain-specific analytical queries are seen in scientific databases, the DBMS

should support efficient processing of those that are frequently used as building blocks for more complex analysis. However, many of such basic analytical queries need super-linear processing time if handled in a straightforward way, as in current scientific databases. In this paper, we report our efforts to design efficient algorithms for a type of query that are extremely important in the analysis of **particle simulation data**.

Particle simulations are computer simulations in which the basic components (e.g., atoms, stars, etc.) of large systems (e.g., molecules, galaxies, etc.) are treated as classical entities that interact for certain duration under postulated empirical forces. For example, molecular simulations (MS) explore relationship between molecular structure, movement and function. These techniques are primarily applicable in modeling of complex chemical and biological systems that are beyond the scope of theoretical models. MS has become an important research tool in material sciences [6], astrophysics [7], biomedical sciences, and biophysics [8], motivated by a wide range of applications. In astrophysics, the N-body simulations are predominantly used to describe large scale celestial structure formation [8]–[11]. Similar to MS in applicability and simulation techniques, the N-body simulation comes with even larger scales in terms of total number of particles simulated.

Results of particle simulations form large datasets of particle configurations. Typically, these configurations store information about the particle types, their coordinates and velocities - the same type of data we have seen in spatial-temporal databases [12]. While snapshots of configurations are interesting, quantitative structural analysis of inter-atomic structures are the mainstream tasks in data analysis. This requires the calculation of statistical properties or functions of particle coordinates [9]. Of special interest to scientists are those quantities that require coordinates of two particles simultaneously. In their brute force form these quantities require  $O(N^2)$  computations for  $N$  particles [8]. In this paper, we focus on one such analytical query: the **Spatial Distance Histogram (SDH)** query, which asks for a histogram of the distances of all pairs of particles in the simulated system.

### A. Problem statement

The SDH problem can be defined as follows: given the coordinates of  $N$  points in space, we are to compute the counts of point-to-point distances that fall into a series of  $l$  ranges in the  $\mathbb{R}$  domain:  $[r_0, r_1), [r_1, r_2), [r_2, r_3), \dots, [r_{l-1}, r_l]$ . A range  $[r_i, r_{i+1})$  in such series is called a *bucket*,

and the span of the range  $r_{i+1} - r_i$  is called the *width* of the bucket. In this paper, we focus our discussions on the case of *standard SDH queries* where all buckets have the same width  $p$  and  $r_0 = 0$ , which gives the following series of buckets:  $[0, p), [p, 2p), \dots, [(l-1)p, lp]$ . Generally, the boundary of the last bucket  $lp$  is set to be the maximum distance of any pair of points in the dataset. Although almost all scientific data analysis only require the computation of standard SDH queries, our solutions can be easily extended to handle histograms with non-uniform bucket width and/or arbitrary values of  $r_0$  and  $r_l$ .<sup>1</sup> The SDH is basically a series of non-negative integers  $\mathbf{h} = (h_1, h_2, \dots, h_l)$  where  $h_i$  ( $0 < i \leq l$ ) is the number of pairs of points whose distances are within the bucket  $[(i-1)p, ip)$ .

### B. Motivation

The SDH is a fundamental tool in the validation and analysis of particle simulation data. It serves as the main building block of a series of critical quantities to describe a physical system. Specifically, SDH is a direct estimation of a continuous statistical distribution function called *radial distribution functions* (RDF) [7], [9], [13]. The RDF is defined as

$$g(r) = \frac{N(r)}{4\pi r^2 \delta r \rho} \quad (1)$$

where  $N(r)$  is the expected number of atoms in the shell between  $r$  and  $r + \delta r$  around any particle,  $\rho$  is the average density of particles in the whole system, and  $4\pi r^2 \delta r$  is the volume of the shell. Since SDH directly provides the value for  $N(r)$ , the RDF can be viewed as a normalized SDH.

The RDF is of great importance in computation of thermodynamic quantities about the simulated system. Some of the important quantities like total pressure,

$$p = \rho kT - \frac{2\pi}{3} \rho^2 \int dr r^3 u'(r) g(r, \rho, T)$$

and energy

$$\frac{E}{NkT} = \frac{3}{2} + \frac{\rho}{2kT} \int dr 4\pi r^2 u(r) g(r, \rho, T)$$

<sup>1</sup>The only complication of non-uniform bucket width is that, given a distance value, we need  $O(\log l)$  time to locate the bucket instead of constant time for equal bucket width.

cannot be calculated without  $g(r)$ . For mono-atomic systems, the RDF can also be directly related to the structure factor of the system [14], via

$$S(k) = 1 + \frac{4\pi\rho}{k} \int_0^\infty (g(r) - 1) r \sin(kr) dr.$$

We skip the definitions of all notations in the above formulae, as the purpose is to show the importance of SDH in particle simulations. To compute SDH in a straightforward way, we have to calculate distances between all pairs of particles and put the distances into bins with a user-specified width, as done in state-of-the-art simulation data analysis software packages [7], [15]. MS or N-body techniques generally consist of large number of particles. For example, the Virgo consortium has accomplished a simulation containing 10 billion particles to study the formation of galaxies and quasars [16]. MS systems also hold up to millions of atoms. This kind of scale prohibits the analysis of large datasets following the brute-force approach. From a database viewpoint, it would be desirable to make SDH a basic query type with the support of scalable algorithms.

Previous work [17], [18] have addressed this problem by developing algorithms that compute exact SDHs with time complexity lower than quadratic. The main idea is to organize the data in a space-partitioning tree and process pairs of tree nodes instead of pairs of particles (thus saving processing time). The tree structure used include  $kd$ -tree in [17] and region quad/oct-tree in our previous work [18], which also proved that the time complexity of such algorithms is  $O(N^{\frac{2d-1}{2d}})$  where  $d \in \{2, 3\}$  is the number of dimensions in the data space. While beating the naive solution in performance, such algorithms' running time for large datasets can still be undesirably long. On the other hand, a SDH with some bounded error can satisfy the needs of users. In fact, there are cases where even a coarse SDH will greatly help the fine-tuning of simulation programs [9]. Generally speaking, the main motivation to process SDHs is to study the statistical distribution of point-to-point distances in the simulated system [9]. Since a histogram by itself is an approximation of the underlying distribution  $g(r)$  (Equation 1), an inaccurate histogram generated from a given dataset will still be useful in a statistical sense. Therefore, in this paper, we focus on approximate algorithms with very high performance and deliver query results with low error rates. In addition to experimental results, we also evaluate the performance/accuracy tradeoffs provided by the proposed algorithms in an analytical way. In short, the running time of our proposed algorithm is only related to the accuracy that needs to

be achieved. In practice, our algorithm achieves excellent performance/accuracy tradeoff – the error rates in query results are very small even when the running time is reasonably short.

### *C. Roadmap of the paper*

We continue this paper by a survey of related work and a list of our contributions in Section II; we introduce the technical background on which our approximate algorithm is built in Section III; we describe the details of a basic approximate algorithm and relevant empirical evaluation in Section IV; then we dedicate Section VI to mathematical analysis of the key mechanisms in our basic algorithm; the results and suggestions of our analytical work is used to develop a new algorithm with improved performance and we introduce and evaluate that algorithm in Section VII; finally, we conclude this paper by Section VIII.

## II. RELATED WORK AND OUR CONTRIBUTIONS

The scientific community has gradually moved from processing large data files towards using database systems for the storage, retrieval, and analysis of large-scale scientific data [2], [19]. Conventional (relational) database systems are designed and optimized toward data and applications from the business world. In recent years, the database community has invested much efforts into constructing database systems that are suitable for handling scientific data. For example, the BDBMS project [3] handles annotation and provenance of biological sequence data; and the PeriScope [5] project is aimed at efficient processing of declarative queries against biological sequences. In addition to that, there are also proposals of new DBMS architectures for scientific data management [20]–[22]. The main challenges and possible solutions of scientific data management are discussed in [1].

Traditionally, molecular simulation data are stored in large files and queries are implemented in standalone programs, as represented by popular simulation/analytics packages [15]. Recent efforts have been dedicated to building simulation data management systems on top of relational databases, as represented by the BioSimGrid [4] and SimDB [23] projects developed for molecular simulations. However, such systems are still in short of efficient query processing strategies. To the best of our knowledge, the computation of SDH in such software packages is done in a brute-force way, which requires  $O(N^2)$  time.

In particle simulations, the computation of (gravitational/electrostatic) force is of similar flavor to the SDH problem. Specifically, the force (or potential) is the sum of all pairwise interactions in the system, thus requires  $O(N^2)$  steps to compute. The simulation community has adopted approximate solutions represented by the Barnes-Hut algorithm that runs on  $O(N \log N)$  time [24] and the Multi-pole algorithm [25] with linear running time. Although all above algorithms use a tree-like data structure to hold the data, they provide little insights on how to solve the SDH problem. The main reason is that these strategies take advantage of two features of force: 1) for any pairwise interaction, its contribution to the force decreases dramatically when particle distance increases; 2) the effects of symmetric interactions cancel out. However, neither features are applicable to SDH computation, in which every pairwise interaction counts and all are equally important. Another method for force computation is based on well-separated pair decomposition (WSPD) [26] and was found to be equivalent to the Barnes-Hut algorithm. A WSPD is a collection of pairs of subsets of the data such that all point-to-point distances are covered by such pairs. The pairs of subsets are also well-separated in that the smallest distance between the smallest balls covering the subsets (with radius  $r$ ) is at least  $sr$  where  $s$  is a user-defined parameter. Although relevant by intuition, the WSPD does not produce fast solution for SDH computation.

Although SDH is an important analytics, there is not much elaboration on efficient SDH algorithms. An earlier work from the data mining community [17] opened the direction of processing SDHs by space-partitioning trees. The core idea is to process all the particles in a tree node as one single entity to take advantage of the non-zero bucket width  $p$ . By this, processing time is saved by avoiding computation of particle-to-particle distances. Our earlier paper [18] proposed a similar algorithm as well as rigorous mathematical analysis (which is not found in [17]) of the algorithm's time complexity. Specifically, in [18], we propose a novel algorithm (named *DM-SDH*) to compute SDH based on a data structure called *density map*, which can be easily implemented by augmenting a Quad-tree index. On contrary to that, the data structure adapted in [17] is the kd-tree. Our mathematical analysis [27] has shown that the algorithm runs on  $\Theta(N^{\frac{3}{2}})$  for two-dimensional data and  $\Theta(N^{\frac{5}{3}})$  for three-dimensional data, respectively. The technical details of such an algorithm will be introduced in Section III.

This paper significantly extends our earlier work [18] by focusing on the approximate algorithms for SDH processing. In particular, we claim the following contributions via this work:

1. We present an approximate SDH processing strategy that is derived from the basic exact algorithm, and this approximate algorithm has constant-time complexity and a provable error bound;
2. We develop a mathematical model to analyze the effects of error compensation that led to high accuracy of our algorithm; and
3. We propose an improved approximate algorithm based on the insights obtained from the above analytical results.

### III. PRELIMINARIES

In this section, we introduce the algorithm we developed in [18] to compute exact SDHs. Techniques and analysis related to this algorithm are the basis for the approximate algorithm we focus on in this paper. In Table I, we list the notations that are used throughout this paper. Note that symbols defined and referenced in a local context are not listed here.

TABLE I  
SYMBOLS AND NOTATIONS.

Symbol	Definition
$p$	width of histogram buckets
$l$	total number of histogram buckets
$\mathbf{h}$	the histogram with elements $h_i$ ( $0 < i \leq l$ )
$N$	total number of particles in data
$i$	an index symbol for any series
$DM_i$	the $I$ -th level density map
$d$	number of dimensions of data
$\epsilon$	error bound for the approximate algorithm
$H$	total level of density maps, i.e., tree height

#### A. Overview of the Density Map-based SDH (DM-SDH) Algorithm

To beat the  $O(N^2)$  time needed by the naive solution, we need to avoid the computation of all particle-to-particle distances. An important observation here is: a histogram bucket always has a non-zero width  $p$ . Given a pair of points, their bucket membership could be determined



if we only know a range that the distance belongs to and this range is contained in a histogram bucket. The central idea of our approach is a conceptual data structure called *density map*. For a 3D space, a density map is essentially a 3D grid that divides the simulated space into cubes of equal volumes. For a 2D space, it consists of squares of equal size. From now on, we use 2D data and grids to elaborate our ideas unless specified otherwise. Note that extending our discussions to 3D data/space would be straightforward. In every cell of the grid, we record the number of particles that are located in the space represented by that cell as well as the four coordinates that determine the exact boundary of the cell in space. The reciprocal of the cell size in a density map is called the *resolution* of the density map. In order to process SDH, we build a series of density maps with different resolutions. We organize the array of density maps in a way such that the resolution of a density map is always doubled as compared to the previous one in the series. Consequently, any cell in a density map is divided into exactly four (eight for a 3D space) disjoint cells in the next density map. A natural way to organize the density maps is to connect all cells in a point region (PR) Quad-tree [28].

The pseudocode of the DM-SDH algorithm can be found in Fig. 1. The core of the algorithm is a procedure named RESOLVETWOCELLS, which is given as inputs a pair of cells  $M_1$  and  $M_2$  on the same density map. In RESOLVETWOCELLS, we first compute the minimum and maximum distances between any particle from  $M_1$  and any one from  $M_2$  (line 1). Obviously, this can be accomplished in constant time given the corner coordinates of two cells stored in the density map. When the minimum and maximum distances between  $M_1$  and  $M_2$  fall into the same histogram bucket  $i$ , we say these two cells are *resolvable* on this density map, and they *resolve* into bucket  $i$ . If this happens, the histogram is updated (lines 2 - 5) by incrementing the count of the specific bucket  $i$  by  $n_1 n_2$  where  $n_1, n_2$  are the particle counts in cells  $M_1$  and  $M_2$ , respectively. If the two cells do not resolve on the current density map, we move to a density map with higher (doubled) resolution and repeat the previous step. However, on this new density map, we try resolving all four partitions of  $M_1$  with all those of  $M_2$  (lines 12 - 16). In other words, there are  $4 \times 4 = 16$  recursive calls to RESOLVETWOCELLS if  $M_1$  and  $M_2$  are not resolvable on the current density map. In another scenario where  $M_1$  and  $M_2$  are not resolvable yet no more density maps are available, we have to calculate the distances of all particles in the non-resolvable cells (lines 6 - 11). The DM-SDH algorithm starts at the first density map  $DM_o$  whose cell diagonal length is smaller than the histogram bucket width  $p$  (line 2). It is easy to

**Algorithm DM-SDH****Inputs:** all data points, density maps built beforehand, and bucket width  $p$ **Output:** an array of counts  $\mathbf{h}$ 

```

1  initialize all elements in  $\mathbf{h}$  to 0
2  find the first density map  $DM_o$  whose cells have diagonal length  $k \leq p$ 
3  for all cells in  $DM_o$ 
4  do  $n \leftarrow$  number of particles in the cell
5      $h_0 \leftarrow h_0 + \frac{1}{2}n(n-1)$ 
6  for any two cells  $M_j$  and  $M_k$  in  $DM_o$ 
7  do RESOLVETWOCELLS ( $M_j, M_k$ )
8  return  $\mathbf{h}$ 

```

Procedure RESOLVETWOCELLS ( $M_1, M_2$ )

```

0  check if  $M_1$  and  $M_2$  are resolvable
1  if  $M_1$  and  $M_2$  are resolvable
2  then  $i \leftarrow$  index of the bucket  $M_1$  and  $M_2$  resolve into
3      $n_1 \leftarrow$  number of particles in  $M_1$ 
4      $n_2 \leftarrow$  number of particles in  $M_2$ 
5      $h_i \leftarrow h_i + n_1 n_2$ 
6  else if  $M_1$  and  $M_2$  are on the last density map (i.e., leaf level of the quad-tree)
7     for each particle A in  $M_1$ 
8     for each particle B in  $M_2$ 
9     do  $f \leftarrow$  distance between A and B
10         $i \leftarrow$  the bucket  $f$  falls into
11         $h_i \leftarrow h_i + 1$ 
12 else
13     $DM' \leftarrow$  next density map with higher resolution
14    for each partition  $M'_1$  of  $M_1$  on  $DM'$ 
15    for each partition  $M'_2$  of  $M_2$  on  $DM'$ 
16    do RESOLVETWOCELLS ( $M'_1, M'_2$ )

```

Fig. 1. The density-map-based SDH algorithm.

see that no pairs of cells are resolvable in density maps with resolution lower than that of  $DM_o$ . Within each cell on  $DM_o$ , we are sure that any intra-cell point-to-point distance is smaller than  $p$  thus all such distances are counted into the first bucket with range  $[0, p)$  (lines 3 - 5). The algorithm proceeds by resolving inter-cell distances (i.e., calling RESOLVETWOCELLS) for all pairs of cells in  $DM_o$  (lines 6 - 7).

Clearly, by only considering atom counts in the density map cells (i.e., quad tree nodes), we are able to process multiple point-to-point distances between two cells in one shot. This translates into significant performance improvements over the brute-force approach.

In DM-SDH, we assume there are a series of density maps built beforehand in the form of a quad-tree. An important implementation detail that is relevant to our approximate algorithm design is the height of the quad tree (i.e., the number of density map levels). Recall that DM-SDH saves time by resolving cells such that we need not to calculate the point-to-point distances one by one. However, when the total point counts in a cell decreases, the time we save by resolving that cell also decreases. Imagine a cell with only 4 or fewer (8 for 3D data/space) data points, it does not give us any benefit in processing SDH to further partition this cell on the next level: the cost of resolving the partitions could be higher than directly retrieving the particles and calculating distances (lines 7 - 11 in RESOLVETWOCELLS). Based on this observation, the total level of density maps  $H$  is set to be

$$H = \left\lceil \log_{2^d} \frac{N}{\beta} \right\rceil + 1 \quad (2)$$

where  $d$  is the number of dimensions,  $2^d$  is the degree of the nodes in the tree (4/8 for 2D/3D data) and  $\beta$  is the average number of particles we desire in each leaf node. In practice, we set  $\beta$  to be slightly greater than 4 in 2D (8 for 3D data) since the CPU cost of resolving two cells is higher than computing the distance between two points.

### B. Performance Analysis of DM-SDH

We have accomplished a rigorous analysis of the performance of DM-SDH and derived its time complexity. The analysis focuses on the quantity of number of point-to-point distances that can be covered in resolved cells. We generate closed-form formulae for such quantities via a geometric modeling approach therefore rigorous analysis of the time complexity becomes possible. While the technical details of the analytical model are complex and can be found in a

recent article [27], it is necessary to sketch the most important (and also most relevant) analytical results here for the purpose of laying out a foundation for the proposed approximate algorithm.

*Theorem 1:* For any given standard SDH query with bucket width  $p$ , let  $DM_o$  be the first density map where the DM-SDH algorithm starts running, and  $\alpha(m)$  be the ratio of non-resolvable pairs of cells on a density map that lies  $m$  levels below  $DM_o$  (i.e., map  $DM_{o+m}$ ) to the total number of cell pairs on that density map. We have

$$\lim_{p \rightarrow 0} \frac{\alpha(m+1)}{\alpha(m)} = \frac{1}{2}.$$

*Proof:* See Section 4 of [27]. ■

What Theorem 1 tells us is: the chance that any pair of cells is not resolvable decreases by half with the density map level increases by one. In other words, for a pair of non-resolvable cells on  $DM_j$  where  $j \geq o$ , among the 16 pairs of subcells on the next level, we expect  $16 \times \frac{1}{2} = 8$  pairs to be resolvable. Our analysis also shows that Theorem 1 not only works well for large  $l$  (i.e., smaller  $p$ , and more meaningful in simulation data analysis), but also **quickly converges even when  $l$  is reasonably small**. Furthermore, the above result is also true for **3D data** (see Section 5.1 of [27]). The importance of Theorem 1 is in that it shows the number of pairs of cells that do not resolve declines exponentially when the algorithm visits more levels of the density map. This is critical in studying the time complexity of DM-SDH.

Given a SDH query with parameter  $p$ , the starting level  $DM_o$  is fixed in DM-SDH. Suppose there are  $I$  non-resolvable pairs of cells on  $DM_o$ . On the next level  $DM_{o+1}$ , total number of cell pairs becomes  $I2^{2d}$ . According to Theorem 1, half of them will be resolved, leaving only  $I2^{2d-1}$  pairs unresolved. On level  $DM_{o+2}$ , the number of non-resolvable pairs of cells becomes  $\frac{I2^{2d-1}2^{2d}}{2} = I2^{2(2d-1)}$ . Thus, the number of calls to resolve cells made by DM-SDH is

$$T_c(N) = I(1 + 2^{2d-1} + 2^{2(2d-1)} + \dots + 2^{n(2d-1)}) = \frac{I[2^{(2d-1)(n+1)} - 1]}{2^{2d-1} - 1} \quad (3)$$

According to Eq. (2), **one more level of density map will be built when data size increases from  $N$  to  $2^d N$** . This gives the following recurrence

$$T_c(2^d N) = \frac{I[2^{(2d-1)(n+2)} - 1]}{2^{2d-1} - 1} = 2^{2d-1} T_c(N) - o(1) \quad (4)$$

which derives

$$T_c(N) = O(N^{\log_{2d} 2^{2d-1}}) = O(N^{\frac{2d-1}{d}}).$$

Given Theorem 1, it is easy to find that the number of distances in the non-resolvable cells also follows the same recurrence as shown in Eq. (4) (see Section 6 of [27]). By that, we conclude the time complexity of the DM-SDH algorithm is  $O(N^{\frac{2d-1}{d}})$ .

#### IV. THE APPROXIMATE DENSITY MAP-BASED SDH ALGORITHM

In this section, we introduce a modified SDH algorithm to give such approximate results to gain better performance in return. Our solution targets at two must-have features of a decent approximate algorithm :1) provable and controllable error bounds such that the users can have an idea on how close the results are to the fact; and 2) analysis of costs to reach (below) a given error bound, which guides desired performance/correctness tradeoffs.

In the DM-SDH algorithm, we have to : 1) keep resolving cells till we reach the lowest level of the tree; 2) calculate point-to-point distances when we cannot resolve two cells on the leaf level of the tree. Our idea for approximate SDH processing is: *stop at a certain tree level and totally skip all distance calculations if we are sure that the number of distances in the unvisited cell pairs fall below some error tolerance threshold*. We name the new algorithm as ADM-SDH (that stands for Approximate Density Map-based SDH), and it can be easily implemented by modifying the DM-SDH algorithm. In particular, we stop the recursive calls to RESOLVETWOCELLS after  $m$  levels. The critical problem, however, is how to determine the value of  $m$  given a user-specified error tolerance bound  $\epsilon$ . In this paper, we use the following metric to quantify the errors

$$e = \frac{\sum_i |h_i - h'_i|}{\sum_i h_i}$$

where for any bucket  $i$ ,  $h_i$  is the accurate count and  $h'_i$  the count given by the approximate algorithm.

For any given density map  $DM_{o+m}$  and total number of buckets  $l$ , our analytical model (Theorem 1) gives the percentage of non-resolvable cell pairs  $\alpha(m)$ . Furthermore, due to the existence of a closed-form formula (see Section 4.4 of [27]),  $\alpha(m)$  can be efficiently computed. We list some values of  $1 - \alpha(m)$ , the percentage of *resolvable* cell pairs, in Table II. Given a user-specified error bound  $\epsilon$ , we can find the appropriate levels of density maps to visit such that the unvisited cell pairs only contain less than  $\epsilon \frac{N(N-1)}{2}$  distances. For example, for a SDH query with 128 buckets and error bound of  $\epsilon = 3\%$ , we get  $m = 5$  by consulting the table. This

TABLE II

EXPECTED PERCENTAGE OF PAIRS OF CELLS THAT CAN BE RESOLVED UNDER DIFFERENT LEVELS OF DENSITY MAPS AND TOTAL NUMBER OF HISTOGRAM BUCKETS. COMPUTED WITH MATHEMATICA 6.0.

Map	Total Number of Buckets				
levels	2	8	32	128	256
1	50.6565	52.5131	52.6167	52.6225	52.6227
2	74.8985	76.2390	76.3078	76.3112	76.3114
3	87.3542	88.1171	88.1539	88.1556	88.1557
4	93.6550	94.0582	94.0777	94.0778	94.0778
5	96.8222	97.0290	97.0285	97.0389	97.0389
6	98.4098	98.5145	98.5198	98.5195	98.5195
7	99.2046	99.2572	99.2596	99.2597	99.2597
8	99.6022	99.6286	99.6298	99.6299	99.6299
9	99.8011	99.8143	99.8149	99.8149	99.8149
10	99.9005	99.9072	99.9075	99.9075	99.9075

means, to ensure the 3% error bound, we only need to visit five levels of the tree (excluding the starting level  $DM_o$ ), and no distance calculation is needed. Table II serves as an excellent validation of Theorem 1:  $\alpha(m)$  almost exactly halves itself when  $m$  increases by 1, even when  $l$  is as small as 2. Since the numbers on the first row (i.e., values for  $1 - \alpha(1)$ ) are also close to 0.5, the correct choice of  $m$  for the guaranteed error rate  $\epsilon$  is

$$m = \lg \frac{1}{\epsilon}.$$

The cost of the approximate algorithm only involves resolving cells on the  $m + 1$  levels of density maps. From formula (4), we obtain the time complexity of the new algorithm

$$T_c(N) \approx I 2^{(2d-1)m} = I 2^{(2d-1) \lg \frac{1}{\epsilon}} = I \left( \frac{1}{\epsilon} \right)^{2d-1} \quad (5)$$

where  $I$  is the number of cell pairs on the starting density map  $DM_o$ , and it is solely determined by the query parameter  $p$ . Apparently, the running time of this algorithm is not related to the input size  $N$ .

#### A. Heuristic Distribution of Distance Counts

Now let us discuss how to deal with those non-resolvable cells after visiting  $m + 1$  levels on the tree. In giving the error bounds in our approximate algorithm, we are conservative in assuming

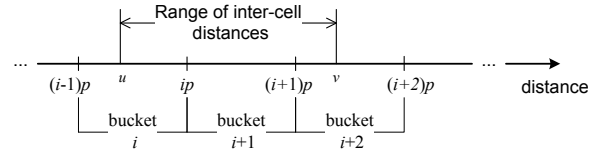


Fig. 2. Distance range of two resolvable cells overlap with three buckets.

the distances in all the unresolved cells will be placed into the wrong bucket. In fact, this almost will never happen because we can distribute the distance counts in the unvisited cells to the histogram buckets heuristically and some of them will be done correctly. Consider two non-resolvable cells in a density map with particle counts  $n_1$  and  $n_2$  (i.e., total number of  $n_1 n_2$  distances between them), respectively. We know their minimum and maximum distances  $u$  and  $v$  (these are calculated beforehand in our attempt to resolve them) fall into multiple buckets. Fig. 2 shows an example that spans three buckets. Using this example, we describe the following heuristics to distribute the  $n_1 n_2$  total distance counts into the relevant buckets. These heuristics are ordered in their expected correctness.

1. Put all  $n_1 n_2$  distance counts into one bucket that is predetermined (e.g., always putting the counts to the leftmost bucket); We name this heuristic as SKEW;
2. Evenly distribute the distance counts into the three buckets that  $[u, v]$  overlaps, i.e., each bucket gets  $\frac{1}{3} n_1 n_2$ ; this heuristic is named EVEN;
3. Distribute the distance counts based on the overlaps between range  $[u, v]$  and the buckets. In Fig. 2, the distances put into buckets  $i$ ,  $i + 1$ , and  $i + 2$  are  $n_1 n_2 \frac{ip - u}{v - u}$ ,  $n_1 n_2 \frac{p}{v - r}$ , and  $n_1 n_2 \frac{v - (i + 1)p}{v - u}$ , respectively. Apparently, by adapting this approach, we assume the (statistical) distribution of the point-to-point distances between the two cells is uniform.

This heuristic is called PROP (short for *proportional*).

The assumption of uniform distance distribution in PROP is obviously an oversimplification. In [29], we briefly mentioned a 4th heuristic: if we know the spatial distribution of particles within individual cells, we can generate the statistical distribution of the distances either analytically or via simulations, and put the  $n_1 n_2$  distances to involved buckets based on this distribution. This solution involves very non-trivial statistical inference of the particle spatial distribution and is beyond the scope of this paper.

Note that all above methods require only constant time to compute a solution for two cells. Therefore, the time complexity of ADM-SDH is not affected no matter which heuristic is used.

## V. EMPIRICAL EVALUATION OF ADM-SDH

We have implemented the ADM-SDH algorithm using the C programming language and tested it with various synthetic/real datasets. The experiments are run at an Apple Mac Pro workstation with two dual-core 2.66GHz Intel Xeon CPUs, and 8GB of physical memory. The operating system is OS X 10.5 Leopard. In these experiments, we set the program to stop after visiting different levels of density maps and distribute the distances using the three heuristics (Section IV). We then compare the approximate histogram with those generated by regular DM-SDH. We use various synthetic and real data sets in our experiments. The synthetic data are generated from: (1) uniform distributions to simulate a system with particles evenly distributed in space; and (2) Zipf distribution with order 1 to introduce skewness to data spatial distribution. The real datasets are extracted from a molecular simulation of biomembrane structures (Fig. 3). The data size in such experiments range from 50,000 to 12,800,000.

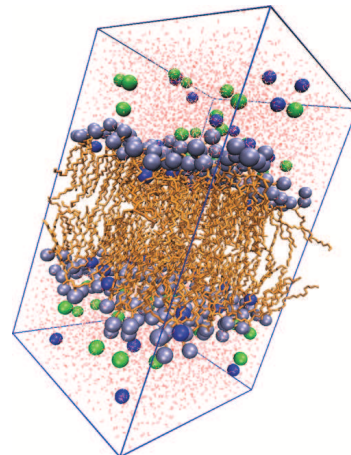


Fig. 3. The simulated hydrated dipalmitoylphosphatidylcholine bilayer system. We can see two layers of hydrophilic head groups (with higher atom density) connected to hydrophobic tails (lower atom density) are surrounded by water molecules (red dots).

Fig. 4 shows the running time of ADM-SDH under one single  $p$  value of 2500.0. Note that the ‘Exact’ line shows the results of the basic DM-SDH algorithm whose running time obviously increases polynomially with  $N$  at a slope of about 1.5. First, we can easily conclude that the running time after the tree construction stage does not change with the increase of dataset size (Fig. 4(a)). The only exception is when  $m$  is 5 – the running time increases when  $N$  is small and then stays as a constant afterwards. This is because the algorithm has less than 5 levels to visit in a bushy tree resulted from small  $N$  values. When  $N$  is large enough, running time no longer changes with the increase of  $N$ . In Fig. 4(b), we plot the total running time which includes the time for quad-tree construction. Under small  $m$  values, the tree construction time is a dominating factor since it increases with data size  $N$  (i.e.,  $O(N \log N)$ ). However, when  $m > 3$ , the shape



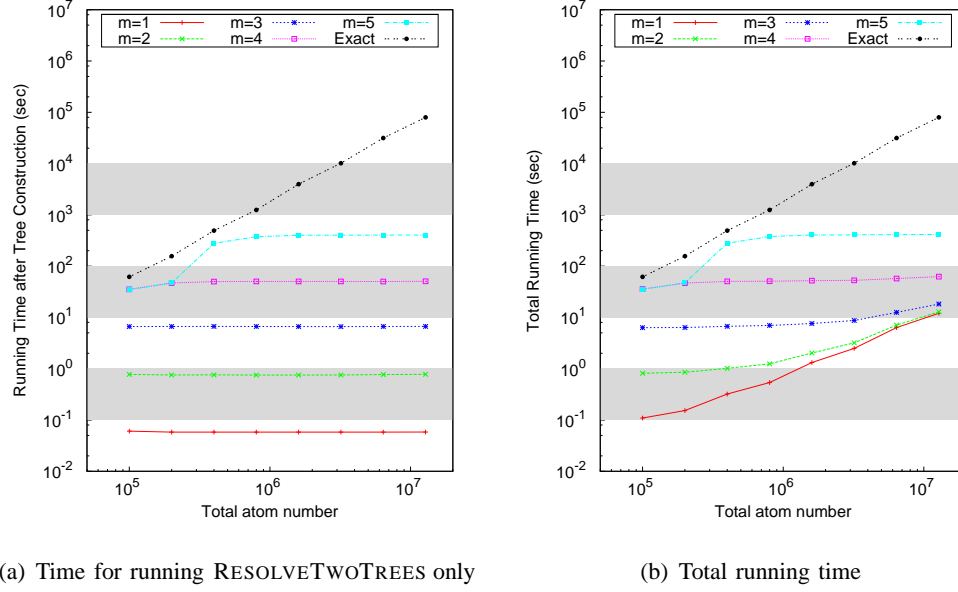


Fig. 4. Efficiency of the ADM-SDH algorithm.

of the curve does not change much as compared to those in Fig. 4(a), indicating the time for running RESOLVETWOTREES dominates.

We observed surprising results on the accuracy of ADM-SDH. In Fig. 5, we plot the error rates observed in experiments with three different datasets and three heuristics mentioned in Section IV-A. First, it is obvious that less error was observed when  $m$  increases. The exciting fact is that, in almost all experiments, the error rate is lower than 10% – even for the cases of  $m = 1$ ! These are much lower than the error bounds we get from Table II. The correctness of heuristic SKEW is significantly lower than that of EVEN, and that of EVEN lower than PROP, as expected. Heuristic PROP achieves very low error rates even in scenarios with small  $m$  values. For most experiments, the data size  $N$  does not seem to affect the error rate of the algorithm. The above trends are observed in all three datasets.

### A. Discussions

At this point, we can conclude that the ADM-SDH algorithm is an elegant solution to the SDH computation problem. According to our experiments, extremely low error rates can be obtained even we only visit as few as one level of density map, leading to a very efficient algorithm yet with high accuracy in practice. It is clearly shown that the required running time

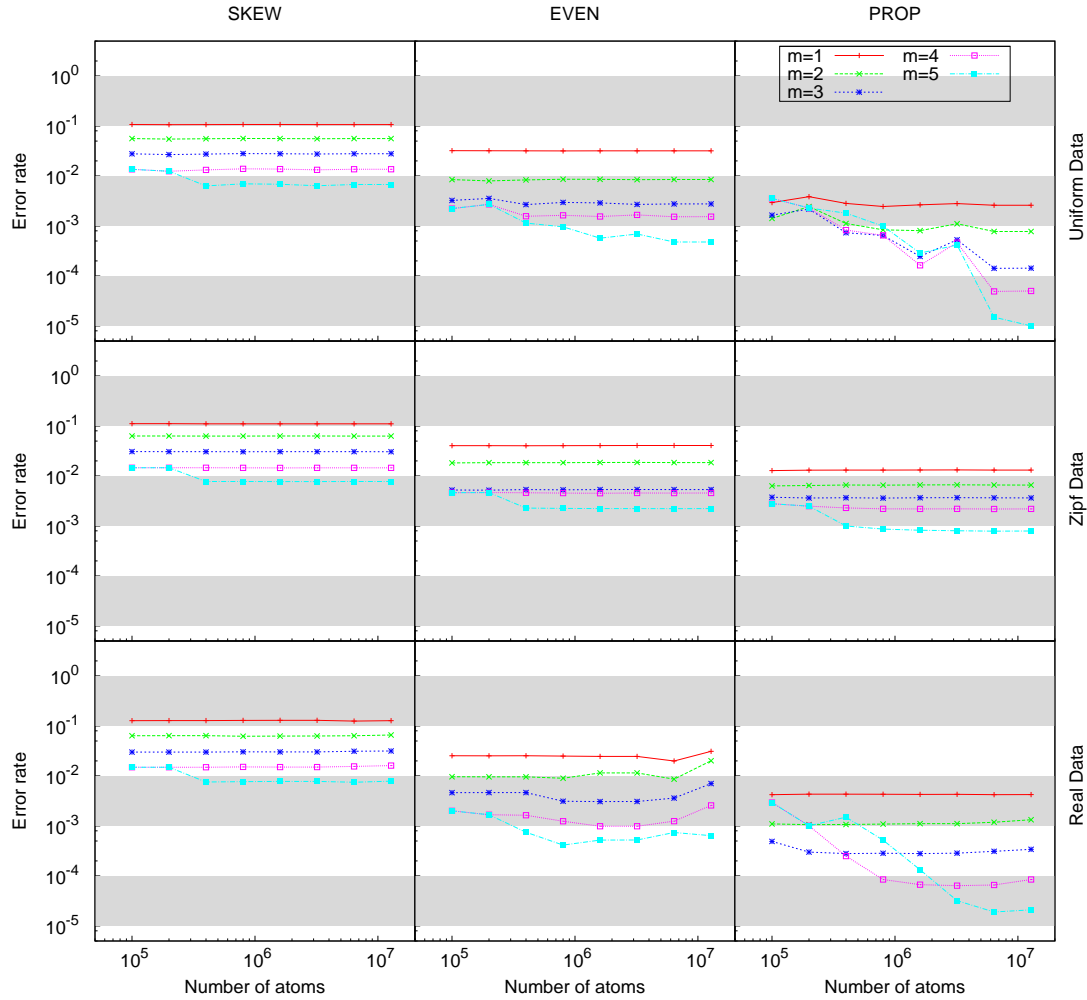


Fig. 5. Accuracy of the ADM-SDH algorithm.

for ADM-SDH grows nicely with the data size  $N$  (i.e., only when  $m$  is of a small value does the tree construction time dominate).

The errors rates achieved by ADM-SDH approximate algorithm shown by current experiments are much lower than what we expected from our basic analysis. For example, Table II predicts an error rate of around 48% for the case of  $m = 1$ , yet the error we observed for  $m = 1$  in our experiments is no more than 10%. With the PROP heuristic, this value can be as low as 0.5%. Our explanation for such low error rates is: in an individual operation to distribute the distance counts heuristically, we could have rendered a large error by putting too many counts into a bucket (e.g., bucket  $i$  in Fig. 2) than needed. But the effects of this mistake could be

(partially) canceled out by another distribution operation, in which too few counts are put into bucket  $i$ . Note that the total error in a bucket is calculated after all operations are processed, thus it reflects the net effects of all positive and negative errors from individual operations. We call this phenomenon *error compensation*.

While more experiments under different scenarios are obviously needed, investigations from an analytical viewpoint are necessary. From the above facts, we understand that the bound given by Table II is loose. The real error bound should be described as

$$\epsilon = \epsilon' \epsilon'' \quad (6)$$

where  $\epsilon'$  is the percentage of resolved distances given by Table II and  $\epsilon''$  is the error rate created by the heuristics via error compensation. In the following section, we develop an analytical model to study how error compensation dramatically boosts accuracy of the algorithm.

## VI. PERFORMANCE ANALYSIS OF ADM-SDH

It is difficult to obtain a tight error bound for ADM-SDH due to the fact that the error is related to data distribution. In this paper, we develop an analytical framework that achieves qualitative analysis of the behaviors of ADM-SDH, with a focus on the generation of errors. Throughout the analysis, we assume uniform spatial distribution of particles and we consider only one level in the density map (i.e.  $m = 1$ ). At the start level (and the only level we visit), the side length of a cell is  $\sqrt{2}p/2$ . First, we are going to analyze the SKEW method.

### A. The distribution of Two Cells' Distance

We study two cells A and B on a density map, with cell A's row number denoted as  $i$  and column number as  $j$ , and cell B's row number as  $k$  and column number as  $l$ . We further denote the minimum distance between A and B as  $u$ , and the maximum distance as  $v$ . We propose the following lemma:

*Lemma 1:* The range  $[u, v]$  overlaps with at most three buckets in the SDH. In other words,  $p \leq v - u \leq 2p$ .

The proof of Lema 1 can be found in Appendix A. By Lemma 1, we can easily see that  $v$  must fall into one of the two buckets with ranges  $[\lfloor \frac{u}{p} \rfloor p + p, \lfloor \frac{u}{p} \rfloor p + 2p)$  and  $[\lfloor \frac{u}{p} \rfloor p + 2p, \lfloor \frac{u}{p} \rfloor p + 3p)$ .

Suppose the distance between points from the two cells follow a cumulative distribution function  $F$  over the range  $[u, v]$ , then the probabilities of a distance falling into the relevant bucket can be found in Table III.

TABLE III

THE BUCKETS INVOLVED IN THE DISTRIBUTION OF DISTANCES FROM TWO NON-RESOLVABLE CELLS.

No	Bucket Range	Cumulative Probabilities
1	$[\lfloor \frac{u}{p} \rfloor p, \lfloor \frac{u}{p} \rfloor p + p)$	$F(\lfloor \frac{u}{p} \rfloor p + p)$
2	$[\lfloor \frac{u}{p} \rfloor p + p, \lfloor \frac{u}{p} \rfloor p + 2p)$	$F(\lfloor \frac{u}{p} \rfloor p + 2p) - F(\lfloor \frac{u}{p} \rfloor p + p)$
3	$[\lfloor \frac{u}{p} \rfloor p + 2p, \lfloor \frac{u}{p} \rfloor p + 3p)$	$1 - F(\lfloor \frac{u}{p} \rfloor p + 2p)$

### B. Compensating the distance counts in the SKEW method

As mentioned earlier, an important mechanism that leads to low error rate in our algorithm is that the errors made by one distribution operation can be compensated by those of another. We can use the SKEW heuristic as an example to study this. In SKEW, all distance counts are put into one bucket, say, the one with the smallest bucket index. In other words, The distance counts in all three buckets (Table III) are put into bucket with range  $[\lfloor \frac{u}{p} \rfloor p, \lfloor \frac{u}{p} \rfloor p + p)$ . The error would be large if we only consider this single distribution operation – by denoting the error as  $e$ , we have  $e = 1 - F(\lfloor \frac{u}{p} \rfloor p + p)$  for the bucket  $[\lfloor \frac{u}{p} \rfloor p, \lfloor \frac{u}{p} \rfloor p + p)$ . The error  $e$  here is positive, meaning counts in the first bucket are overestimated. However, such errors can be canceled out by other distribution operations that move all distance counts from this bucket into another one. For example, if there exists another distribution operation with minimum distance  $u_1 = u - p$ , it would move some counts that belong to bucket 1 in Table III out, generating a negative error in bucket 1 and thus compensating the positive error mentioned before. Given this, an important goal of our analysis is to *find such compensating distribution operations and study how much errors can be compensated*. We first show that under an ideal situation the error can reach zero.

*Lemma 2:* For any distribution operation with minimum distance  $u$ , if there exists another such operation with minimum distance  $u_1 = u - p$ , the error generated by ADM-SDH using the SKEW approach is zero.

The proof of Lema 2 can be found in Appendix B.

In the following text, we study how the errors can be partially compensated by neighboring pairs of cells.

Without loss of generality, we take any pair of cells in the density map – one with coordinates  $(0, 0)$  as the base cell, and another one with coordinates  $(x, y)$ , such as cells  $A$  and  $B$  in Figure 6. For the convenience of presentation, we remove the constant  $p$  in our future discussions. In other words, distances are now defined with  $p$  as the basic unit. For example, the coordinate  $(x, y)$  means that this cell's location is of a distance  $xp$  away from the base cell horizontally and  $yp$  vertically. The minimum distance between the above two cells can be therefore written as (with  $p$  as unit and the side of a cell as  $\frac{\sqrt{2}}{2}p$ ):

$$u = \sqrt{\frac{x^2}{2} + \frac{y^2}{2}} \quad (7)$$

The critical observation that leads to the success of our analysis is obtained by studying another cell located at  $(x - 1, y - 1)$ , i.e., cell  $B'$  in Figure 6. Its minimum distance to the cell  $A$  at  $(0, 0)$  is  $u_1 = \sqrt{\frac{(x-1)^2}{2} + \frac{(y-1)^2}{2}}$ . Let us denote the quantity  $u - u_1$  as  $\Delta$ . We have

$$\begin{aligned} \Delta &= u - u_1 = \frac{(u - u_1)(u + u_1)}{u + u_1} = \frac{u^2 - u_1^2}{u + u_1} \\ &= \frac{\frac{x^2}{2} + \frac{y^2}{2} - \frac{(x-1)^2}{2} - \frac{(y-1)^2}{2}}{u + u_1} = \frac{x + y - 1}{u + u_1} \approx \frac{x + y - 1}{2u} \end{aligned} \quad (8)$$

Suppose  $x \geq y$  and  $z = y/x$ , Equation (8) can be rewritten as follows

$$\begin{aligned} \Delta &\approx \frac{x + y - 1}{2u} = \frac{\frac{x+y-1}{x}}{2\sqrt{\frac{x^2}{2} + \frac{y^2}{2}}} = \frac{1 + \frac{y}{x} - \frac{1}{x}}{\sqrt{4(\frac{x^2}{2x^2} + \frac{y^2}{2x^2})}} \\ &= \frac{1 + z - \frac{1}{x}}{\sqrt{2 + 2z^2}} \approx \frac{1 + z}{\sqrt{2 + 2z^2}} \end{aligned} \quad (9)$$

Although  $x$  and  $y$  are integers, we can treat  $z$  as a continuous variable due to the existence of a large number of possible values of  $x$  and  $y$  in a density map with many cells. Since  $d\Delta/dz > 0$ , we conclude that  $\Delta$  increases with  $z$ . Two boundary cases are: 1) when  $y = 0$  we have  $z = 0$  and  $\Delta = \frac{\sqrt{2}}{2}$ ; 2) when  $y = x$ , we have  $z = 1$  and  $\Delta = 1$ .

According to Lemma 2, the error of the SKEW method,  $e_{skew}$ , can be noted by the difference between one and  $\Delta$ ,  $e_{skew} = 1 - \Delta$ . This error,  $e_{skewed}$ , ranges from 0 to  $1 - \frac{\sqrt{2}}{2}$ , since  $\Delta$  ranges

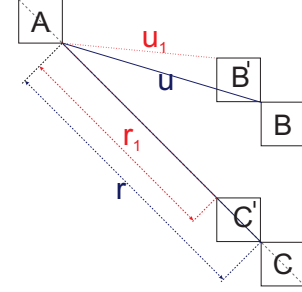


Fig. 6. Pairs of cells that lead to partial error compensation in the SKEW approach.

from  $\frac{\sqrt{2}}{2}$  to 1. The compensating process in distance counts is shown in Figure 6. As mentioned before,  $C$  and  $C'$  are examples of two cells for which the difference of minimum distances to cell  $A$  is 1 and the cells  $B$  and  $B'$  are two cells for which the difference of minimum distances is different from one (less than one in this case).

Figure 7 represents the triangular density distributions of three different minimum distances,  $u$ ,  $u_1$  and  $u - 1$ . In our analysis we have used normal distribution in stead of triangular, based on the approximation shown in appendix C.  $u$  is the minimum distance between the base cell and the cell which is located at  $(x, y)$  and  $u_1$  is the minimum distance between the base cell and the cell which is located at  $(x - 1, y - 1)$ . In our analysis we define  $G$  as  $\lfloor u \rfloor$ , and we also define  $H = G + 1$  and  $I = H + 1$ . The triangles of  $STU$ ,  $AEW$  and  $BCF$  are the density distributions of the minimum distances  $u$ ,  $u_1$  and  $u - 1$ , respectively. The line  $H'S'D'$  is the symmetry line of  $WFH$  with respect to the vertical line which passes through the point  $D$ .

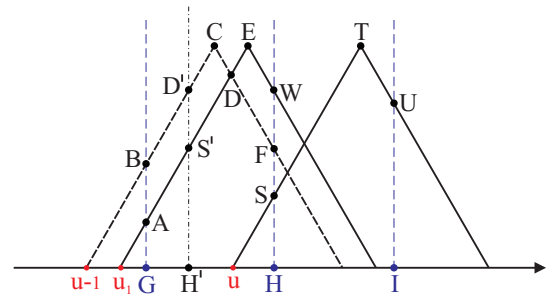


Fig. 7. Compensating the distance counts

As we know, if  $u_1 = u - 1$ , the error of distance counts produced by the period between  $H$  and  $I$  can be compensated by that produced by the period between  $G$  and  $H$ . In other words, according to SKEW, when we count the number of distances between the cell at coordinates  $(0, 0)$  and the cell located at  $(x, y)$ , the number of distances between  $H$  and  $I$  is added and leads to a positive error. When we count the number of distances between the base cell and the cell which is located at  $(x - 1, y - 1)$ , the number of distances between  $G$  and  $H$  is missed and leads to a negative error. If the distribution is the same, the area of  $GBCFHH$  is the same as the area of  $HSTUI$ .

If  $u_1 \neq u - 1$ , there is difference between the area of  $GAEWH$  and the area of  $HSTUI$ . In other words, there is difference between the area of  $GAEWH$  and the area of  $GBCFHH$ . This difference can be computed by the area of  $ABD'S'$  and that represents the error imposed by the difference between

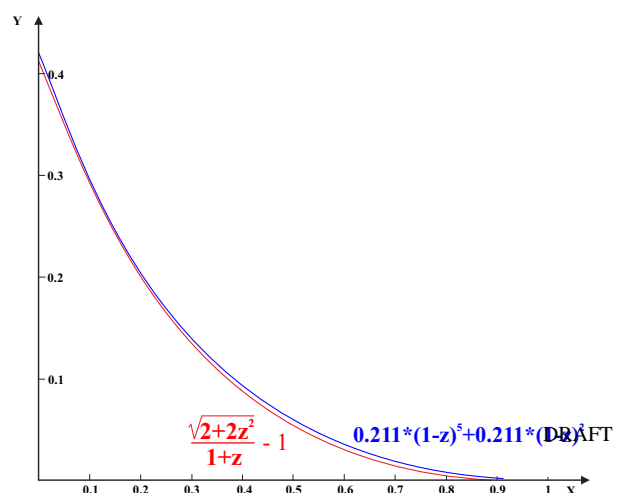


Fig. 8. Approximation of  $\frac{1}{1+z} - 1$

$u_1$  and  $u - 1$  which we denote as  $e_{u_1, u-1}$ . We

know that  $CE = u_1 - u + 1$  and  $GH' <$

$u - u_1 = \Delta$ . Considering the fact that the

area of each of the three triangles in Figure 7

is one, the height of each of these triangles is  $\frac{2}{v-u}$  (since the base is  $v - u$ ). Therefore, the ratio  $\frac{AB}{CE}$  has the following value (for more detailed explanation, please refer to Appendix D):

$$\frac{AB}{CE} = \frac{\frac{2}{v-u}}{\frac{v-u}{2}} = \frac{4}{(v-u)^2} \quad (10)$$

Furthermore, the length of  $AB$  is:

$$AB = \frac{4}{(v-u)^2} * CE = \frac{4(u_1 - u + 1)}{(v-u)^2} = \frac{4(1-\Delta)}{(v-u)^2} \quad (11)$$

The area of  $ABD'S'$ , thus the error  $e_{u_1, u-1}$ , can be computed as follows:

$$e_{u_1, u-1} = AB * GH' < \frac{4(1-\Delta) * \Delta}{(v-u)^2} \quad (12)$$

$$e_{u_1, u-1} < \frac{(1-\Delta)}{\Delta} = \frac{1}{\Delta} - 1 \quad (13)$$

Therefore the accumulated error,  $e_{z=\overline{0,1}}$  over the range of  $z$ ,  $z = \overline{0,1}$  can be computed by the following equation (considering Eq. (9) for  $\Delta$ )

$$e_{z=\overline{0,1}} = \sum_{z=0}^1 e_{u_1, u-1} = \sum_{z=0}^1 \left( \frac{1}{\Delta} - 1 \right) \approx \sum_{z=0}^1 \left( \frac{\sqrt{2+2z^2}}{1+z} - 1 \right) \quad (14)$$

When  $0 \leq z \leq 1$ , we can approximate  $e_{z=\overline{0,1}}$  as follows (from Figure 8).

$$e_{z=\overline{0,1}} \leq \sum_{z=0}^1 0.211 * (1-z)^5 + 0.211 * (1-z)^2 \quad (15)$$

$$e_{z=\overline{0,1}} = \int_0^1 (0.211 * (1-z)^5 + 0.211 * (1-z)^2) dz = 0.1055 \quad (16)$$

Eq. (16) means that the total error rendered by the SKEW under the assumptions we stated in the beginning of Section VI is less than 10.55%. Due to the assumptions we made, we do not claim this as a rigorous bound. However, it clearly shows that our algorithm is able to produce really good results with low errors by visiting only one level in the density map. And this conclusion builds the foundation of an improved approximate algorithm (Section VII).

One special note here is that Eq. (16) does not cover the cases in which the minimum distance  $u$  falls into the first SDH bucket (i.e.,  $u < p$ ). However, our analysis shows that such cases do not impact the results in Eq. (16) significantly. More details can be found in Appendix E.

## VII. SINGLE LEVEL APPROXIMATE ALGORITHM

Via the performance analysis of ADM-SDH, and looking back to the error bound described in Eq. (6), we concluded that  $\epsilon''$  is very small. Even if we allow  $\epsilon'$  be 100%, meaning no cell resolution is needed, we can still achieve low and controllable error rates in our results. Therefore, in this section, we introduce an improved approximate algorithm, we call *single level approximate algorithm* (SL-SDH), based on such a conclusion. SL-SDH has all the same parts as the original ADM-SDH algorithm with the only difference is that we now *visit only one density map (i.e., one level of the tree) instead of visiting  $m$  levels as in the original approximate algorithm*. SL-SDH improves over ADM-SDH in two important aspects. First, we only need a single  $DM$  which can be built in  $O(N)$  time (instead of the  $O(N \log N)$  time needed to build the quad tree). Second, we reduce the post-tree-construction running time of the algorithm, with little increase of the error, as we only run RESOLVETWOTREES for cells in one density map.

However, the running time of SL-SDH still depends heavily on the bucket width  $p$ . Recall that ADM-SDH starts at the density map  $DM_0$  where the diagonal of a single cell is less than or equal to  $p$ . When  $p$  is small, the number of cells in  $DM_0$  is large, and we have to invoke the RESOLVETWOCELLS procedure more times. To remedy this, we further modify SL-SDH by allowing it to run on a (single) density map above  $DM_0$ , i.e., one with larger cell sizes (and fewer cells). This is based on a hypothesis motivated by our performance analysis of ADM-SDH: the error compensation mechanism we studied will also work for density maps above  $DM_0$ . We know the error is very small for running, RESOLVETWOTREES for those cells in  $DM_0$  - doing the same on higher level density maps should still render reasonable (although higher) error rates. Unfortunately, an analytical study of such errors is very difficult. In the remainder of this section, we empirically evaluate the error and time tradeoff of the final version of the SL-SDH algorithm.



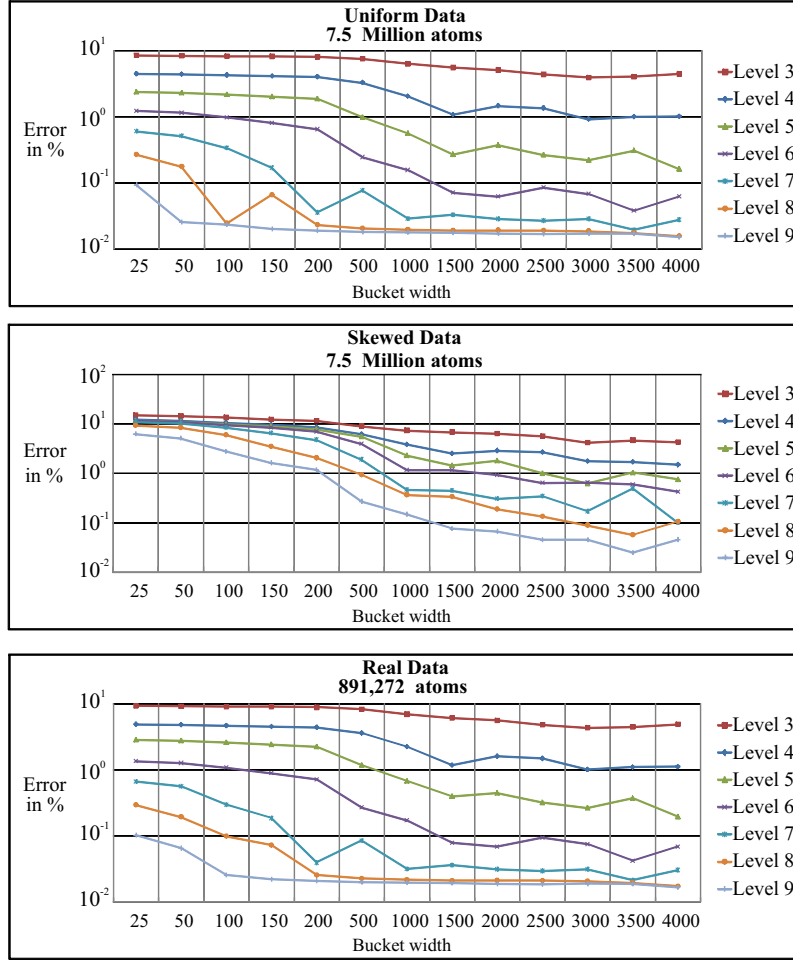


Fig. 9. Accuracy of the SL-SDH algorithm under different bucket width for synthetic (uniform and skewed) and real data.

### A. Experimental Results

We have implemented the SL-SDH algorithm using the C programming language and tested it with various synthetic/real datasets. The experiments are run in the same environment as the experiments for the ADM-SDH in Section V.

Since we know, from our previous experiments, that the PROP heuristic for distributing the distances in non-resolvable cells produces the best results, we have only used that heuristic to show the results of the single level approximate algorithm. We have run the algorithm on two synthetic datasets, one with uniform and another with skewed distribution of atoms, under five different numbers (i.e., 1, 3, 5, 7.5 and 12 million) of atoms. We also ran the algorithm on one real simulation dataset with 891,272 atoms. Figure 9 shows the results from the experiments on uniform and skewed data respectively with 7.5 million atoms and the real data with 891,272

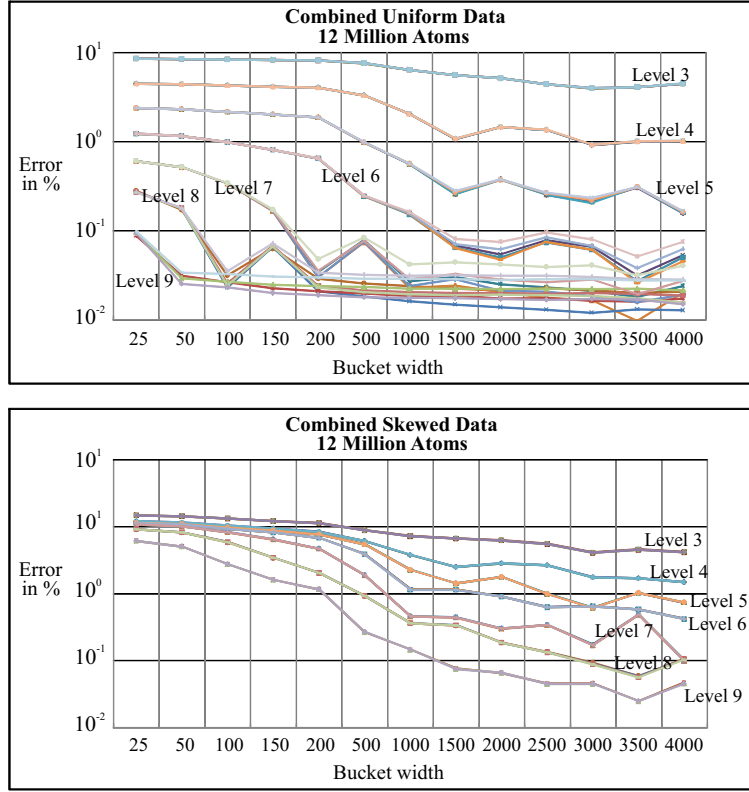


Fig. 10. Accuracy of the SL-SDH algorithm under different bucket width and different atom counts for synthetic data (uniform and skewed).

atoms. From these figures we can see that the error rate decreases when we increase the level in the density map. We also see that the error rate is lower when the bucket width increases.

In Figure 10 we demonstrate the effects of system size  $N$  on the accuracy of SL-SDH. Each line in Figure 10 plots the error rates of SL-SDH when run with a particular level of density map under a particular atom count. We can easily see that the lines for the five different system sizes (of the same density map) are very similar, giving rise to one cluster of lines for each level of density map. The above two figures actually show that the error introduced by the SL-SDH algorithm is not affected by the number of atoms in the dataset  $N$ , but by the level of density map the algorithm works at. On the other hand, the running time of the SL-SDH algorithm was also found to be independent of  $N$ , as shown in Figure 11. The only exceptions are for levels 3 and 4, in which the tree construction time dominates.

The difference between SL-SDH and ADM-SDH can be seen by comparing Figure 11 with Figure 4, and Figure 10 with Figure 5. However, to better understand the accuracy/performance

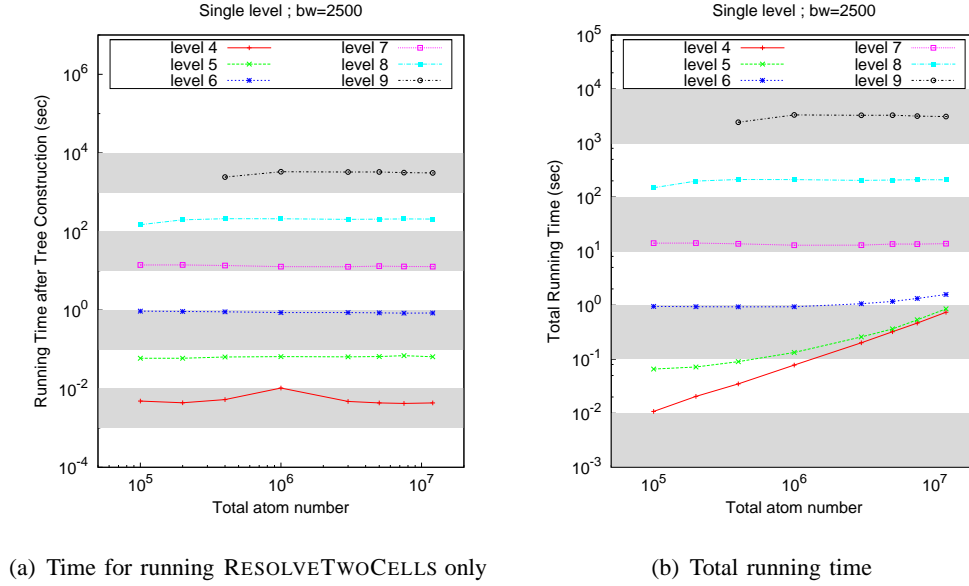


Fig. 11. Efficiency of the SL-SDH algorithm.

tradeoffs provided by the two algorithms, we introduce a new metric **Error Delay Product (EDP)**, which is defined as the product of the error rate and the running time of the algorithm. Obviously, higher EDP means worse accuracy/performance tradeoff. Figures 12, 13 and 14 show the EDP for both algorithms run under three different bucket width (i.e., 100, 500 and 1000). From these figures it is obvious that the SL-SDH algorithm shows better EDP than the regular approximate algorithm. This is especially the case when the bucket width is small – the EDPs of SL-SDH are orders of magnitude lower than those provided by ADM-SDH. The reason for this is that the ADM-SDH algorithm starts at a level determined by the bucket width  $p$ . When  $p$  gets smaller, it has to start at a lower level of the tree with more cells to process. In Figure 14, ADM-SDH shows better EDPs, but the lowest level it reaches is still several times higher than that of SL-SDH. In general, the EDP of SL-SDH decreases with the decrease of density map level, with level 9 always shows the worst tradeoff. This means the accuracy of our approximate algorithm is already high with working on a coarse density map, the improvement of accuracy levels out with further increase of the density map resolution.

In summary, our experimental results convey two important messages. First, the SL-SDH algorithm significantly improves the accuracy/performance tradeoff over ADM-SDH. Such improvements are more obvious under small SDH bucket width. This is very important: the

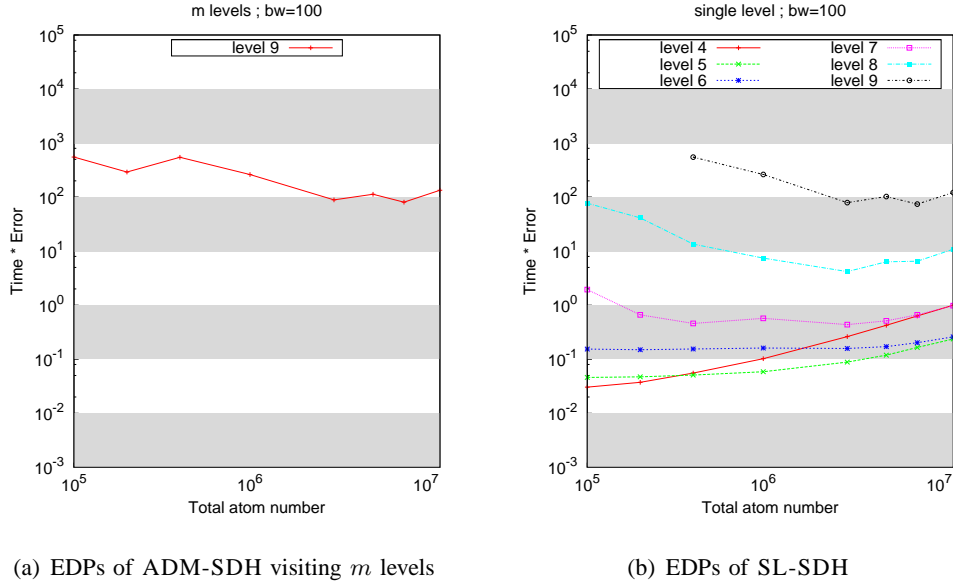


Fig. 12. Accuracy and performance tradeoffs of the ADM-SDH and SL-SDH algorithms under SDH bucket width of 100.

computation of SDHs is generally preferred to be done under smaller  $p$  values as it carries more information about the distribution of the distances. Users can choose the appropriate (single) level among all the density maps to run the algorithm based only on the desired accuracy. On the other hand, we also show that, like ADM-SDH, the running time and the error rate of SL-SDH are not affected by the number of atoms in the data set.

## VIII. CONCLUSIONS AND FUTURE WORK

The main objective of our work is to accomplish efficient computation of SDH, a popular quantity in particle simulations, with guaranteed accuracy. In this paper, we introduce approximate algorithm for SDH processing based on our proviso work developed around a Quadtree-like data structure named *density map*. The experimental results show that our approximate algorithm has very high performance (short running time) while delivering results with astonishingly low error rates. Aside from the experimental results, we also analytically evaluate the performance/accuracy tradeoffs of the algorithm. Such analyses showed that the running time of our algorithm is completely independent of the input size  $N$ , and derived a provable error bound under desired running time. We further developed another mathematical model to perform in-depth study of the mechanism that leads to low error rates of the algorithm. Aside from administering tighter

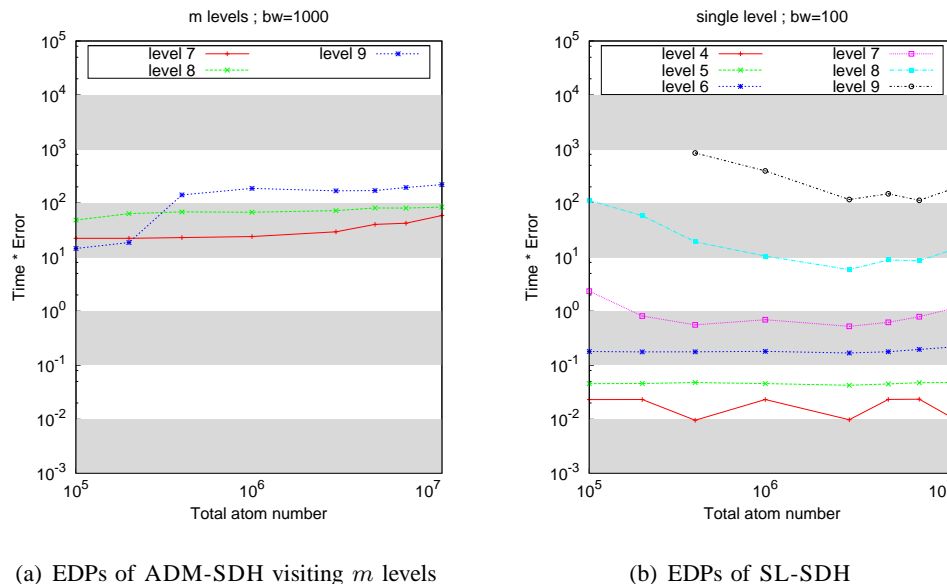


Fig. 13. Accuracy and performance tradeoffs of the ADM-SDH and SL-SDH algorithms under SDH bucket width of 500.

bounds (under some assumptions) on the error of the basic approximate algorithm, our model also gives insights on how the basic algorithm can be improved. Following these insights, a new single level approximate algorithm with improved time/accuracy tradeoff was proposed. Our experimental results supported our analysis. Having these experimental results on hand, one aspect of our future work will be to establish a provable error bound for the new algorithm.

Many times, the molecular simulation systems are observed over certain period of time and SDH computation is required for every frame (time instance) over that period. Therefore, another direction of our on-going work is to efficiently compute the SDHs of consecutive frames by taking advantage of the temporal locality of data points. We can also extend our work to the computation of  $m$ -body correlation functions with  $m > 2$  – a more general form of SDH that involves counting all possible  $m$ -particle tuples.

## APPENDIX A

### Proof of Lemma 1

We study two cells A and B on a density map, with cell A's row number denoted as  $i$  and

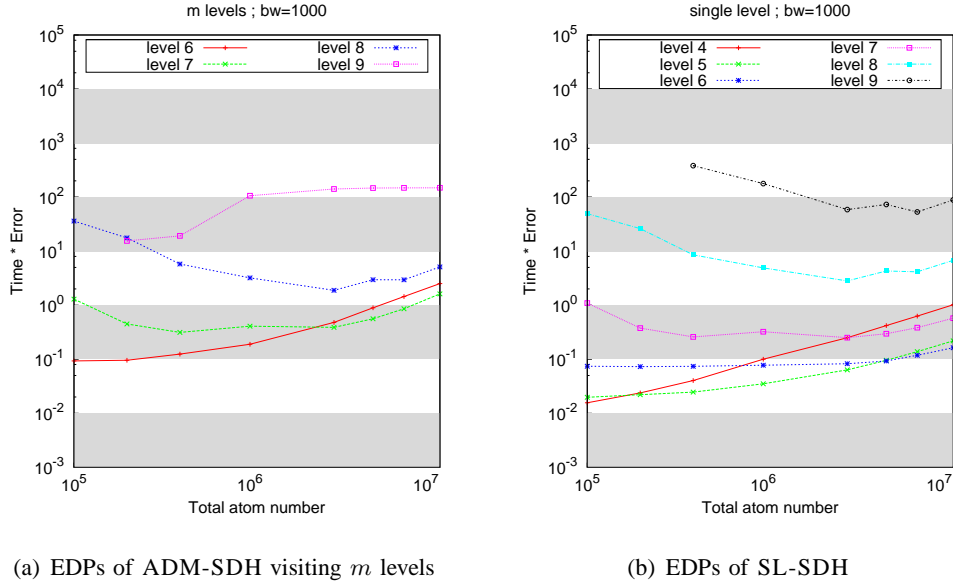


Fig. 14. Accuracy and performance tradeoffs of the ADM-SDH and SL-SDH algorithms under SDH bucket width of 1000.

column number as  $j$ , and cell B's row number as  $k$  and column number as  $l$ . We further denote the minimum distance between A and B as  $u$ , and the maximum distance as  $v$ . Then  $v$  and  $u$  can be written as functions of the two cells' row and column numbers as follows:

**Case 1.** When A and B are in different rows and columns, i.e.,  $i \neq k$  and  $j \neq l$ , we have

$$u = p \sqrt{\frac{(i - k - 1)^2}{2} + \frac{(j - l - 1)^2}{2}} \quad (17)$$

and

$$v = p \sqrt{\frac{(i - k + 1)^2}{2} + \frac{(j - l + 1)^2}{2}} \quad (18)$$

**Case 2.** When Cell A is located in the same row (but not in the same column) as Cell B, i.e.,  $i = k$ ,  $r$  and  $v$  can be rewritten as follows

$$u = p \frac{|j - l - 1| \sqrt{2}}{2} \quad (19)$$

$$v = p \sqrt{\frac{(j - l + 1)^2}{2} + \frac{1}{2}} \quad (20)$$

**Case 2.** When Cell A is located in the same column (but not in the same row) as Cell B, i.e.,  $j = l$ ,  $r$  and  $v$  can be rewritten as follows

$$u = p \frac{|i - k - 1| \sqrt{2}}{2} \quad (21)$$

$$v = p\sqrt{\frac{(i-k+1)^2}{2} + \frac{1}{2}} \quad (22)$$

We consider the following two scenarios to accomplish the proof.

**Case 1.** Cell A is located on the same row or column as Cell B. By equations (19), (20), (21) and (22), we can easily get

$$v - u > 2p * \sqrt{2}/2 > p$$

and

$$v - u = p \frac{4\sqrt{2}r + 5}{2v + 2u} < p \frac{4\sqrt{2}r + 5}{4u} = \left( \sqrt{2} + \frac{5}{4u} \right) p$$

When  $u \geq 2\sqrt{2}p$ , we can easily see that  $\left( \sqrt{2} + \frac{5}{4u} \right) p < 2p$  and thus  $v - u < 2p$ . If  $u < 2\sqrt{2}$ , we have to study the value of quantity  $v - u$  case by case. Fortunately,  $u$  can only be of a series of discrete values. We enumerate such cases as follows:

When  $u = 0$  (i.e., the cells are adjacent to each other), we have  $v - r = v = p\sqrt{10}/2$ ;

When  $u = \frac{\sqrt{2}}{2}p$ ,  $v = \sqrt{5}p$ ,  $v - u \approx 1.646p < 2p$ .

When  $u = \sqrt{2}p$ ,  $v = \frac{\sqrt{34}}{2}p$ ,  $v - u \approx 1.52p < 2p$ .

When  $u = \frac{3\sqrt{2}}{2}p$ ,  $v = \sqrt{13}p$ ,  $v - u \approx 1.5p < 2p$ .

When  $u = 2\sqrt{2}p$ ,  $v = \frac{\sqrt{74}}{2}p$ ,  $v - u \approx 1.5p < 2p$ .

**Case 2.** When cell A and cell B are located in different rows and columns (as shown in Fig. 15), we can divide  $u$  and  $v$  into two line segments, respectively. First, we have  $u = CB + BD$ , and  $v = AB + BE$ . Since  $AC + CB \geq AB$  and  $BD + DE \geq BE$ , we easily get  $AC + DE + CD \geq AE$ , which is equivalent to  $AC + DE + u \geq v$ . Due to  $AC = DE = p$ , we conclude  $v - u \leq 2p$ .

On the other hand, since  $u$  is the minimum distance between the two cells (i.e.,  $FG \geq u$ ), we have  $v - u = AE - CD \geq AF + GE \geq \sqrt{2}p > p$ .

The results on the above two scenarios conclude the proof.

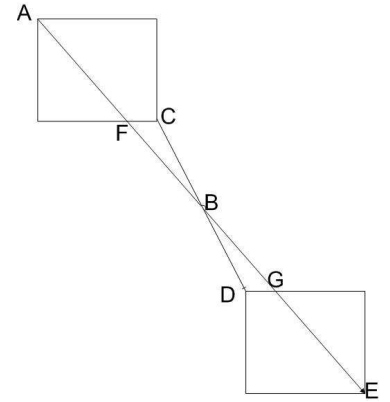


Fig. 15. The minimum and the maximum distance of two cells.

## APPENDIX B

### Proof of Lemma 2

According to Table III, for any distribution operation, the error to the first SDH bucket (denoted as bucket  $i$ ) it involves is  $1 - F(\lfloor \frac{u}{p} \rfloor p + p)$ , and this error is positive (i.e., overestimation). Suppose that there is another distribution with minimum distance  $u_1 = u - p$ , then this operation generates a negative error  $F(\lfloor \frac{u}{p} \rfloor p + 2p) - F(\lfloor \frac{u}{p} \rfloor p + p)$  to bucket  $i$ . For the same reason, a third distribution with minimum distance  $u_2 = u - 2p$  generates a negative error of  $1 - F(\lfloor \frac{u}{p} \rfloor p + 2p)$ . It is easy to see that the combined error (by putting all relevant negative and positive errors together) to bucket  $i$  is 0. An example of two cells that contribute to each other's error compensation in the aforementioned way can be seen in Figure 6. Namely, the cells  $C$  and  $C'$  when we compute the minimum distances  $AC$  and  $AC'$ .

Unfortunately, the above condition of the existence of a  $u_1$  value that equals  $u - p$  cannot be satisfied for all pairs of cells. From Lemma 2, however, we can easily see that the error is strongly related to the quantity  $u$ .

## APPENDIX C

### Approximation of Triangular to Normal Distribution

*Lemma 3:* If  $X$  and  $Y$  are independent random variables uniformly distributed on  $(a, b)$  and  $(c, c + b - a)$ , and  $c \geq a$ , then  $Y - X$  is a triangular random variable and can be regarded as a normal random with the relative error less than 0.1.

*Proof:*

The probability density of  $X$  is

$$f(x) = \frac{1}{b-a}, a < x < b \quad (23)$$

and the probability density of  $Y$  is

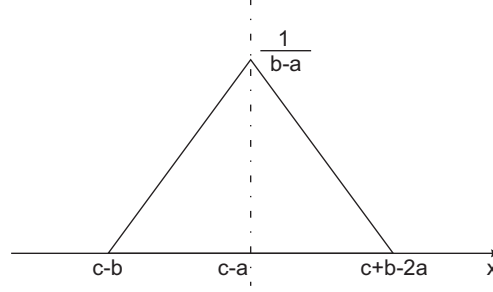
$$g(y) = \frac{1}{b-a}, c < y < c + b - a \quad (24)$$

There are two cases to be considered. One is when  $c$  is equal to  $a$  and the other is when  $c$  is greater than  $a$ .

*Case1:*  $c$  is equal to  $a$

The probability density of  $Y - X$  can be calculated as follows




 Fig. 16. the distribution of  $Y-X$ 

$$f_{Y-X}(z) = \int_c^{c+b-a} f(y-z)g(y)dy \quad (25)$$

When  $0 > z > a - b$ , the probability density of  $Y - X$  can be computed as follows

$$f_{Y-X}(z) = \frac{b-a+z}{(b-a)^2} \quad (26)$$

When  $0 < z < b - a$ , the probability density of  $Y - X$  can be computed as follows

$$f_{Y-X}(z) = \frac{b-a-z}{(b-a)^2} \quad (27)$$

*Case2: c is greater than a.*

When  $c - b < z < c - a$ , the probability density of  $Y - X$  can be computed as follows

$$f_{Y-X}(z) = \frac{b-c+z}{(b-a)^2} \quad (28)$$

When  $c - a < z < c + b - a$ , the probability density of  $Y - X$  can be computed as follows

$$f_{Y-X}(z) = \frac{c+b-2a-z}{(b-a)^2} \quad (29)$$

The probability density of  $Y - X$  is shown in Figure 16.

If  $w$  is a normal random variable with parameters  $(c - a, \frac{(b-a)^2}{6})$ , then the probability density of  $w$  can be written as follows.

$$f(w) = \frac{\sqrt{6}e^{-\frac{6(w-c+a)^2}{2(b-a)^2}}}{\sqrt{2\pi}(b-a)} \quad (30)$$

Let  $u = \frac{w-(c-a)}{b-a}$ . Then, the probability density of  $w$  can be rewritten as follows

$$f(w) = \frac{\sqrt{6}e^{-3u^2}}{\sqrt{2\pi}}$$

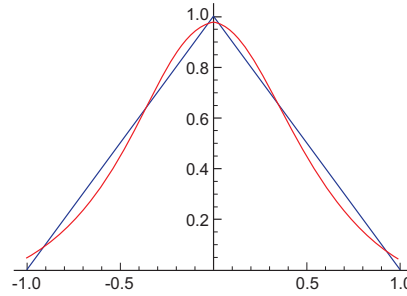


Fig. 17. the relative error between normal and triangular distribution

Let  $v = \frac{z-(c-a)}{b-a}$ . Then, the probability density of  $Y - X$  can be rewritten as follows

$$f_{Y-X}(z) = \text{Min}\{1 + v, 1 - v\}, -1 < v < 1 \quad (31)$$

If  $u = v$ , then, when  $u = \pm 0.02$ ,  $u = \pm 0.41$  or  $u = \pm 0.92$ , the probability density of  $w$  is equal to that of  $Y - X$ . So, if we use the normal distribution to approximate the triangular distribution, the relative error,  $Re$ , which is shown in Figure 17, is given by the following equation:

$$\begin{aligned} Re = & 2 \int_0^{0.02} \left(1 - u - \frac{\sqrt{6}e^{-3u^2}}{\sqrt{2\pi}}\right) du \\ & + 2 \int_{0.02}^{0.41} \left(\frac{\sqrt{6}e^{-3u^2}}{\sqrt{2\pi}} - 1 + u\right) du \\ & + 2 \int_{0.41}^{0.92} \left(1 - u - \frac{\sqrt{6}e^{-3u^2}}{\sqrt{2\pi}}\right) du \\ & + 2 \int_{0.92}^{1.0} \left(\frac{\sqrt{6}e^{-3u^2}}{\sqrt{2\pi}} - 1 + u\right) du \end{aligned} \quad (32)$$

Computing the value for  $Re$ , we get  $Re = 0.093 < 0.1$ . So, we conclude that we can use normal distribution instead of triangular distribution, introducing an error of less than 10%.

■

## APPENDIX D

### Similar Triangles Ratio

Figure 18 represents the two leftmost triangles from Figure 7. These two triangles, triangles  $MNC$  and  $KPE$ , are equivalent with area of one. Their sides  $MC$  and  $KE$  are parallel to each

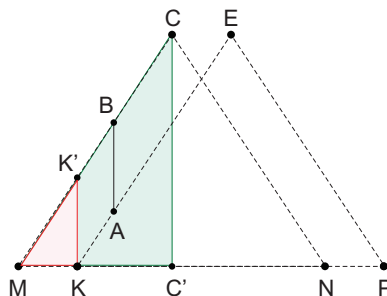


Fig. 18. similar triangles

other. This imposes that  $MK = CE$  and  $KK' = AB$ . Following the values from Figure 7, we note the following:

$MK = u_1 - u + 1 = 1 - \Delta$ ,  $MN = v - u$ ,  $MC' = \frac{MN}{2} = \frac{v-u}{2}$ ,  $CC' = \frac{2}{MN} = \frac{2}{v-u}$  (from the formula of the area of the triangle:  $1 = \frac{1}{2}MN * CC'$ ).

Now, let's look at triangles  $MKK'$  (the red triangle) and triangle  $MC'C$  (the green one). These two triangles are similar right triangles. Following the properties of similar triangles we get the following:

$$\begin{aligned}
\frac{K'K}{MK} &= \frac{CC'}{MC'} = \frac{\frac{2}{v-u}}{\frac{v-u}{2}} = \frac{4}{(v-u)^2} \\
K'K &= \frac{4}{(v-u)^2} * MK \\
AB = K'K &= \frac{4}{(v-u)^2} * MK = \frac{4}{(v-u)^2} * (1 - \Delta)
\end{aligned} \tag{33}$$

## APPENDIX E

### Boundary situation in error compensation analysis

When the minimum distance  $u$  is within the first bucket, i.e.  $u \leq p$ , the error produced by SKEW cannot be compensated in the way described in Section VI. The reason being, there are no buckets to the left of the first bucket, therefore the positive errors in it cannot be compensated.

The average error caused by this, denoted as  $e_{u \leq 1}$ , can be computed as follows

$$e_{u \leq 1} = \frac{0.3N^2}{2N(N-1) + N^2} \approx 0.1 \quad (34)$$

Therefore, by considering the Eqs. (34) and (16), the average error over all buckets of our algorithm based on SKEW can be computed as follows:

$$e_{final} = \lim_{n \rightarrow \infty} \frac{e_{u \leq 1} * \tau + e_{z=\overline{0,1}} * (n - \tau)}{n} = \lim_{n \rightarrow \infty} \frac{0.1 * \tau + 0.1055 * (n - \tau)}{n} = 0.1055 \quad (35)$$

where  $n$  is the total number of distances,  $\tau$  is the number of the distances in the first bucket.

The reason for the second step of the above equation is because  $\tau$  can be computed as:

$$\frac{\tau}{n} = P((x_1 - x_2)^2 + (y_1 - y_2)^2 < 1) \quad (36)$$

Since  $x_1, x_2, y_1$  and  $y_2$  are independent variables with uniform distribution, the above probability can be shown to be very small:

$$P = \frac{1}{2n^4} - \frac{4}{3n^3} + \frac{3\pi - 4}{3n^2} \quad (37)$$

## APPENDIX F

### Distance distribution of particles in two cells A and B

*Lemma 4:* If the points in cell  $A$  and cell  $B$  are uniformly distributed, then the distances between the two cells' points can be described as a Noncentral chi-square distribution and the distribution has no relationship with the number of points in cell  $A$  or cell  $B$ .

*Proof:*

Suppose a point  $i$  in cell  $A$  with coordinates  $x_i$  and  $y_i$  and a point  $j$  in cell  $B$  with coordinates  $x_j$  and  $y_j$ . As we know, if  $x_i$  and  $x_j$  are independent random variables uniformly distributed on  $(a, b)$  and  $(c, c + b - a)$  respectively, ( $c \geq a$ ) then  $x_j - x_i$  follows a triangular distribution. According to the Appendix C, the triangular distribution can be approximated with normal, introducing an error of not more than 10%. Knowing this,  $x_j - x_i$  can be regarded as a normal random variable with parameters  $(c - a, (b - a)^2/6)$ . Similarly,  $y_j - y_i$  can be regarded as a normal random variable with parameters  $(c' - a', (b - a)^2/6)$ . Since  $(b - a)^2/6$  is a constant,

which is noted as  $\sigma^2$ , we can write the following equation for the distance  $l$  between the two points  $i$  and  $j$ :

$$\frac{l^2}{\sigma^2} = \frac{(x_j - x_i)^2}{\sigma^2} + \frac{(y_j - y_i)^2}{\sigma^2} \quad (38)$$

As it is known, the right hand side of the above equation is a Noncentral chi-square distribution. This means that the distances between the two cells' points can be described as a Noncentral chi-square distribution with the parameters  $(2, \lambda)$ , where  $\lambda$  can be defined as follows:

$$\lambda = \frac{(c - a)^2}{\sigma^2} + \frac{(c' - a')^2}{\sigma^2} \quad (39)$$

where  $c - a$  and  $c' - a'$  are the means of the normal distribution.

Since the parameters of Noncentral chi-square distribution have no relationship with the number of points in cell  $A$  or cell  $B$ , the shape of the distribution of the distances is not influenced by the points count in the cells.

So, our conclusion is that the square of the distance between any two points from two cells follows (can be approximated to) a Noncentral chi-square distribution. Since the pdf of a Noncentral chi-square distribution has a closed form which involves an infinite sum, the pdf of the square root of Noncentral chi-square distribution can be easily obtained through Jacobian transformation. So  $E(d)$  can be easily worked out. However, such a closed-form is for an approximated distribution of  $d$  using the technique presented in Appendix C. It is our belief, based on the thorough work we have done on this matter, that to get a closed form for the distribution of the distances between the points of the cells is really tough and challenging, if not impossible to solve, problem.

■

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## REFERENCES

- [1] J. Gray, D. Liu, M. Nieto-Santisteban, A. Szalay, D. DeWitt, and G. Heber, "Scientific Data Management in the Coming Decade," *SIGMOD Record*, vol. 34, no. 4, pp. 34–41, December 2005.
- [2] A. S. Szalay, J. Gray, A. Thakar, P. Z. Kunszt, T. Malik, J. Raddick, C. Stoughton, and J. vandenBerg, "The SDSS Skyserver: Public Access to the Sloan Digital Sky Server Data," in *Proceedings of International Conference on Management of Data (SIGMOD)*, 2002, pp. 570–581.
- [3] M. Y. Eltabakh, M. Ouzzani, and W. G. Aref, "BDBMS - A Database Management System for Biological Data," in *Proceedings of the 3rd Biennial Conference on Innovative Data Systems Research (CIDR)*, 2007, pp. 196–206.
- [4] M. H. Ng, S. Johnston, B. Wu, S. E. Murdock, K. Tai, H. Fangohr, S. J. Cox, J. W. Essex, M. S. P. Sansom, and P. Jeffreys, "BioSimGrid: Grid-enabled Biomolecular Simulation Data Storage and Analysis," *Future Generation Computer Systems*, vol. 22, no. 6, pp. 657–664, June 2006.
- [5] J. M. Patel, "The Role of Declarative Querying in Bioinformatics," *OMICS: A Journal of Integrative Biology*, vol. 7, no. 1, pp. 89–91, 2003.
- [6] S. Klasky, B. Ludaescher, and M. Parashar, "The Center for Plasma Edge Simulation Workflow Requirements," in *EEE Workshop on Workflow and Data Flow for Scientific Applications (SciFlow'06)*, 1991, pp. 73–73.
- [7] J. L. Stark and F. Murtagh, *Astronomical Image and Data Analysis*. Springer, 2002.
- [8] M. P. Allen and D. J. Tildesley, *Computer Simulations of Liquids*. Clarendon Press, Oxford, 1987.
- [9] D. Frenkel and B. Smit, *Understanding Molecular Simulation: From Algorithm to Applications*, ser. Computational Science Series. Academic Press, 2002, vol. 1.
- [10] J. M. Haile, *Molecular Dynamics Simulation: Elementary Methods*. Wiley, New York, 1992.
- [11] D. P. Landau and K. Binder, *A Guide to Monte Carlo Simulation in Statistical Physics*. Cambridge University Press, Cambridge, 2000.
- [12] P. K. Agarwal, L. Arge, and J. Erikson, "Indexing Moving Objects," in *Proceedings of International Conference on Principles of Database Systems (PODS)*, 2000, pp. 175–186.
- [13] M. Bamdad, S. Alavi, B. Najafi, and E. Keshavarzi, "A new expression for radial distribution function and infinite shear modulus of lennard-jones fluids," *Chem. Phys.*, vol. 325, pp. 554–562, 2006.
- [14] A. Filipponi, "The radial distribution function probed by X-ray absorption spectroscopy," *J. Phys.: Condens. Matter*, vol. 6, pp. 8415–8427, 1994.
- [15] B. Hess, C. Kutzner, D. van der Spoel, and E. Lindahl, "GROMACS 4: Algorithms for Highly Efficient, Load-Balanced, and Scalable Molecular Simulation," *Journal of Chemical Theory and Computation*, vol. 4, no. 3, pp. 435–447, March 2008.
- [16] V. Springel, S. D. M. White, A. Jenkins, C. S. Frenk, N. Yoshida, L. Gao, J. Navarro, R. Thacker, D. Croton, J. Helly, J. A. Peacock, S. Cole, P. Thomas, H. Couchman, A. Evrard, J. Colberg, and F. Pearce, "Simulations of the Formation, Evolution and Clustering of Galaxies and Quasars," *Nature*, vol. 435, pp. 629–636, June 2005.
- [17] A. G. Gray and A. W. Moore, "N-body problems in statistical learning," in *Advances in Neural Information Processing Systems (NIPS)*. MIT Press, 2000, pp. 521–527.
- [18] Y.-C. Tu, S. Chen, and S. Pandit, "Computing Distance Histograms Efficiently in Scientific Databases," in *Proceedings of International Conference on Data Engineering (ICDE)*, March 2009, pp. 796–807.
- [19] M. Arya, W. F. Cody, C. Faloutsos, J. Richardson, and A. Toya, "QBISM: Extending a DBMS to Support 3D Medical Images," in *ICDE*, 1994, pp. 314–325.

- [20] M. Stonebraker, S. Madden, D. J. Abadi, S. Harizopoulos, N. Hachem, and P. Helland, "The End of an Architectural Era (It's Time for a Complete Rewrite)," in *VLDB*, 2007, pp. 1150–1160.
- [21] B. Howe, D. Maier, and L. Bright, "Smoothing the ROI Curve for Scientific Data Management Applications," in *CIDR*, 2007, pp. 185–195.
- [22] P. G. Brown, "Overview of scidb: large scale array storage, processing and analysis," in *SIGMOD Conference*, 2010, pp. 963–968.
- [23] M. Feig, M. Abdullah, L. Johnsson, and B. M. Pettitt, "Large Scale Distributed Data Repository: Design of a Molecular Dynamics Trajectory Database," *Future Generation Computer Systems*, vol. 16, no. 1, pp. 101–110, January 1999.
- [24] J. Barnes and P. Hut, "A Hierarchical  $O(N \log N)$  Force-Calculation Algorithm," *Nature*, vol. 324, no. 4, pp. 446–449, 1986.
- [25] L. Greengard and V. Rokhlin, "A Fast Algorithm for Particle Simulations ," *Journal of Computational Physics*, vol. 135, no. 12, pp. 280–292, 1987.
- [26] P. B. Callahan and S. R. Kosaraju, "A decomposition of multidimensional point sets with applications to k-nearest-neighbors and n-body potential fields," *Journal of ACM*, vol. 42, no. 1, pp. 67–90, 1995.
- [27] S. Chen, Y.-C. Tu, and Y. Xia, "Performance analysis of a dual-tree algorithm for computing spatial distance histograms," *The VLDB Journal*, vol. 20, 2011.
- [28] J. A. Orenstein, "Multidimensional Tries used for Associative Searching," *Information Processing Letters*, vol. 14, no. 4, pp. 150–157, 1982.
- [29] Y.-C. Tu, S. Chen, and S. Pandit, "Computing Distance Histograms Efficiently in Scientific Databases," in *Proceedings of International Conference on Data Engineering (ICDE)*, March 2009, pp. 796–807.